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Keywords Top physics · Charge asymmetry

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1 Introduction

The top quark is the heaviest elementary particle so far observed. With a mass close to the electroweak scale it may play a special role in physics beyond the Standard Model (SM). Its pair production at hadron colliders allows a test of quantum chromodynamics (QCD) at high energies.

This paper describes the measurement of the charge asymmetry A_C , defined as [1, 2]:

$$A_C = \frac{N(\Delta|y| > 0) - N(\Delta|y| < 0)}{N(\Delta|y| > 0) + N(\Delta|y| < 0)}, \quad (1)$$

where $\Delta|y| \equiv |y_t| - |y_{\bar{t}}|$ is the difference between the absolute values of the top and antitop rapidities ($|y_t|$ and

$|y_{\bar{t}}|$) and N is the number of events with $\Delta|y|$ positive or negative.

Although $t\bar{t}$ production at hadron colliders is predicted to be symmetric under the exchange of t and \bar{t} at leading order, at next-to-leading order (NLO) the process $q\bar{q} \rightarrow t\bar{t}g$ exhibits an asymmetry in the differential distributions of the top and antitop, due to interference between initial and final state gluon emission. The $q\bar{q} \rightarrow t\bar{t}$ process also possesses an asymmetry due to the interference between the Born and box diagrams. Similarly, the $gg \rightarrow t\bar{t}q$ process is asymmetric due to interference between amplitudes which have a relative sign difference under the exchange of t and \bar{t} . The production of $t\bar{t}$ pairs by gluon-gluon fusion, $gg \rightarrow t\bar{t}$, on the other hand, is symmetric.

In $p\bar{p}$ collisions at the Tevatron, where top pairs are predominantly produced by quark-antiquark annihilation, perturbative QCD predicts that the top quark will be preferentially emitted in the direction of the incoming quark and the antitop in the direction of the incoming antiquark [3]. Consequently, the charge asymmetry is measured as a forward-backward asymmetry, A_{FB} . Recent measurements of A_{FB} by the CDF and D0 Collaborations [4–7] show a $2\text{--}3\sigma$ excess over the SM expectations enhancing interest in scrutinising the $t\bar{t}$ asymmetry. For $t\bar{t}$ invariant mass, $m_{t\bar{t}}$, greater than 450 GeV, the CDF experiment measures an asymmetry in the $t\bar{t}$ rest frame which is 3.4σ above the SM prediction [6]. Several new physics models have been proposed to explain the excess observed at CDF and D0 [1, 8–17]. Different models predict different asymmetries as a function of $m_{t\bar{t}}$ [18].

In pp collisions at the LHC, the dominant mechanism for $t\bar{t}$ production is expected to be the gluon-gluon fusion process, while $t\bar{t}$ production via $q\bar{q}$ or qg is small. Since the initial state is symmetric, the forward-

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backward asymmetry is no longer a useful observable. However, due to the asymmetry in the production via $q\bar{q}$ and qg , QCD predicts at the LHC a small excess of centrally produced antitop quarks while top quarks are produced, on average, at higher absolute rapidities. This can be understood by the fact that for $t\bar{t}$ production via $q\bar{q}$ annihilation the valence quark carries, on average, a larger momentum fraction than the anti-quark from the sea. With top quarks preferentially emitted in the direction of the initial quarks in the $t\bar{t}$ rest frame, the boost into the laboratory frame drives the top mainly in the forward or backward directions, while antitops are preferentially retained in the central region. If new physics is responsible for the Tevatron A_{FB} excess, the charge asymmetry measured at the LHC is a natural place to look for it.

In this paper, the measurement of the charge asymmetry A_C is performed using candidate $t\bar{t}$ events selected in the lepton+jets channel. In this channel, the SM decay of the $t\bar{t}$ pair to $W^+bW^-\bar{b}$ results in a single electron or muon from one of the W boson decays and four jets, two from the second W boson decay and two from the b - and \bar{b} -quarks. To allow comparisons with theory calculations, the measured $\Delta|y|$ distribution is unfolded to account for acceptance and detector effects. An inclusive measurement, and measurements of A_C in two ranges of $t\bar{t}$ invariant mass, are presented. An inclusive measurement of this asymmetry with an equivalent observable has been recently reported by the CMS collaboration [19].

2 The ATLAS detector

The ATLAS detector [20] at the LHC covers nearly the entire solid angle¹ around the collision point. It consists of an inner tracking detector surrounded by a thin superconducting solenoid, electromagnetic and hadronic calorimeters, and an external muon spectrometer incorporating three large superconducting toroid magnet assemblies.

The inner-detector system is immersed in a 2 T axial magnetic field and provides charged particle tracking in the range $|\eta| < 2.5$. The high-granularity silicon pixel detector covers the vertex region and provides typically

¹ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point (IP) in the centre of the detector and the z -axis along the beam pipe. The x -axis points from the IP to the centre of the LHC ring, and the y axis points upward. Cylindrical coordinates (r, ϕ) are used in the transverse plane, ϕ being the azimuthal angle around the beam pipe. The pseudorapidity is defined in terms of the polar angle θ as $\eta = -\ln \tan(\theta/2)$. Transverse momentum and energy are defined as $p_T = p \sin \theta$ and $E_T = E \sin \theta$, respectively.

three measurements per track, followed by the silicon microstrip tracker (SCT) which provides four measurements from eight strip layers. These silicon detectors are complemented by the transition radiation tracker (TRT), which enables extended track reconstruction up to $|\eta| = 2.0$. In giving typically more than 30 straw-tube measurements per track, the TRT improves the inner detector momentum resolution, and also provides electron identification information.

The calorimeter system covers the pseudorapidity range $|\eta| < 4.9$. Within the region $|\eta| < 3.2$, electromagnetic calorimetry is provided by barrel and endcap lead-liquid argon (LAr) electromagnetic calorimeters, with an additional thin LAr presampler covering $|\eta| < 1.8$ to correct for energy loss in material upstream of the calorimeters. Hadronic calorimetry is provided by the steel/scintillating-tile calorimeter, segmented into three barrel structures within $|\eta| < 1.7$, and two copper/LAr hadronic endcap calorimeters. The solid angle coverage is completed with forward copper/LAr and tungsten/LAr calorimeter modules optimised for electromagnetic and hadronic measurements respectively.

The muon spectrometer comprises separate trigger and high-precision tracking chambers measuring the deflection of muons in a magnetic field with a bending integral from 2 to 8 Tm in the central region, generated by three superconducting air-core toroids. The precision chamber system covers the region $|\eta| < 2.7$ with three layers of monitored drift tubes, complemented by cathode strip chambers in the forward region, where the background is highest. The muon trigger system covers the range $|\eta| < 2.4$ with resistive plate chambers in the barrel, and thin gap chambers in the endcap regions.

A three-level trigger system is used to select interesting events. The level-1 trigger is implemented in hardware and uses a subset of detector information to reduce the event rate to a design value of at most 75 kHz. This is followed by two software-based trigger levels, level-2 and the event filter, which together reduce the event rate to about 300 Hz.

3 Data and Monte Carlo samples

Data from LHC pp collisions collected by the ATLAS detector between March and June 2011 are used in the analysis, corresponding to an integrated luminosity of 1.04 fb^{-1} .

Simulated top pair events are generated using the MC@NLO [21] Monte Carlo (MC) generator with the NLO parton density function (PDF) set CTEQ6.6 [22]. Parton showering and the underlying event are modelled using HERWIG [23] and JIMMY [24], respectively. This $t\bar{t}$ sample is normalised to a cross section

of 165 pb, obtained with the latest theoretical computation, which approximates the next-to-next-to leading order prediction [25]. Single top events are also generated using MC@NLO while the production of W/Z bosons in association with jets is simulated using the ALPGEN generator [26] interfaced to HERWIG and JIMMY with CTEQ6.1 [27]. Diboson events (WW , WZ , ZZ) are generated using HERWIG with MRST2007lomod [28].

All Monte Carlo simulation samples are generated with multiple pp interactions per bunch crossing (pile-up). These simulated events are re-weighted so that the distribution of the number of interactions per crossing in simulation matches that in the data. The samples are then processed through the GEANT4 [29] simulation [30] of the ATLAS detector and the standard reconstruction software.

4 Event selection

4.1 Physics object selection

Reconstructing top quark pair events in the detector requires electrons, muons, jets and missing momentum to be simultaneously measured. Electron candidates are defined as energy deposits in the electromagnetic calorimeter associated with a well-measured track. Identification criteria based on shower shape variables, track quality, and information from the transition radiation tracker are applied to electron candidates [31]. All candidates are required to have $p_T > 25$ GeV and $|\eta_{\text{cluster}}| < 2.47$, where η_{cluster} is the pseudorapidity of the electromagnetic calorimeter cluster associated with the electron. Candidates in the calorimeter transition region $1.37 < |\eta_{\text{cluster}}| < 1.52$ are excluded.

Muon candidates are reconstructed from track segments in different layers of the muon chambers. These segments are combined starting from the outermost layer, with a procedure that takes material effects into account, and matched with tracks found in the inner detector. The candidates are then refitted using the complete track information from both detector systems, and are required to satisfy $p_T > 20$ GeV and $|\eta| < 2.5$.

Jets are reconstructed with the anti- k_t algorithm, with a distance parameter of 0.4 [32], starting from clusters of energy in adjacent calorimeter cells at the electromagnetic (EM) scale. The jet energy is corrected to the hadronic scale using p_T - and η -dependent correction factors obtained from simulation and validated with data [33]. Jet quality criteria are applied to identify jets not associated to in-time real energy deposits in the calorimeters caused by various sources (calorimeter

noise, non-collision beam-related background, cosmic-ray induced showers).

The missing transverse momentum (E_T^{miss}) is reconstructed from clusters of energy calibrated at the EM scale and corrected according to the energy scale of the associated physics object [34]. Contributions from muons are included using their momentum measured from the tracking and muon spectrometer systems. The remaining clusters not associated with the high p_T objects are also included in the missing transverse momentum.

Muons within $\Delta R = 0.4$ of a jet axis² and with $p_T > 20$ GeV are removed in order to reduce the contamination caused by muons from hadron decays. Subsequently, jets within $\Delta R = 0.2$ of an electron candidate are removed to avoid double counting electrons as jets.

Isolation criteria are applied to both electron and muon candidates to reduce the backgrounds from hadrons mimicking lepton signatures and backgrounds from heavy flavour decays inside jets. For electrons, the total energy in a cone of $\Delta R = 0.2$ around the electron candidate must not exceed 3.5 GeV, after correcting for energy deposits from pile-up and for the energy associated with the electron. For muons, the sum of track transverse momenta for all tracks with $p_T > 1$ GeV and the total energy deposited in a cone of $\Delta R = 0.3$ around the muon are both required to be less than 4 GeV ignoring the contribution of the muon p_T .

Reconstructing top quark pair events is facilitated by the ability to tag jets from the hadronisation of b -quarks. For this purpose, two b -tagging algorithms are used and their results are combined to extract a tagging decision for each jet. One b -tagger exploits the topology of b - and c -hadron weak decays inside the jet. A Kalman filter [35] is used to find a common line on which the primary vertex and the b - and c -hadron decay vertices lie, as well as their position on this line, giving an approximate flight path for the b - and c -hadrons. The discrimination between b -, c - and light quark jets is based on a likelihood using the masses, momenta, flight-length significances, and track multiplicities of the reconstructed vertices as inputs. To further increase the flavour discrimination power, a second b -tagger is run which does not attempt to directly reconstruct decay vertices. Instead, this second tagger uses the transverse and the longitudinal impact parameter significances of each track within the jet to determine a likelihood that the jet originates from a b -quark. The results of both taggers are combined using a neural network to determine a single discriminant variable which is used to make tagging decisions. The combined tagger operating

² $\Delta R = \sqrt{\Delta\phi^2 + \Delta\eta^2}$, where $\Delta\phi$ and $\Delta\eta$ are the separation in azimuthal angle and pseudorapidity, respectively.

point chosen for the present analysis corresponds to a 70% tagging efficiency for b -jets in simulated $t\bar{t}$ events while light flavor jets are suppressed by approximately a factor of 100.

4.2 Selection of $t\bar{t}$ candidates

The $t\bar{t}$ final state in the lepton+jets channel is characterised by an isolated lepton (electron or muon) with relatively high p_T , missing transverse momentum arising from the neutrino from the leptonic W decay, two b -quark jets and two light quark jets from the hadronic W decay. To select events with this topology, the appropriate single-electron or single-muon trigger is required to have fired (with thresholds at 20 and 18 GeV respectively). The events are also required to contain one and only one reconstructed lepton with $p_T > 25$ GeV for electrons and $p_T > 20$ GeV for muons. To reject multijet background in the muon channel, $E_T^{\text{miss}} > 20$ GeV and $E_T^{\text{miss}} + m_T(W) > 60$ GeV are required³. In the electron channel more stringent cuts on E_T^{miss} and $m_T(W)$ are required because of the higher level of multijet background, i.e. $E_T^{\text{miss}} > 35$ GeV and $m_T(W) > 25$ GeV. Events are required to have at least four jets with $p_T > 25$ GeV and $|\eta| < 2.5$. These requirements define the ‘pretag’ selection. The ‘tagged’ selection requires, in addition, at least one of the jets with $p_T > 25$ GeV and $|\eta| < 2.5$ to be b -tagged.

5 Background determination

5.1 Multijet background

The method used for evaluating the multijet background with fake leptons⁴ in both the electron and muon channels is the so-called ‘Matrix Method’. This relies on defining loose and tight lepton samples [36] and measuring the fractions of real (ϵ_{real}) and fake (ϵ_{fake}) loose leptons that are selected as tight leptons. The fraction ϵ_{real} is measured using data control samples of Z boson decays to two leptons, while ϵ_{fake} is measured from data control regions defined separately for the electron and muon channels, where the contribution of fake leptons is dominant.

For the muon channel, the loose data sample is defined by removing the isolation requirements in the de-

³Here $m_T(W)$ is the W -boson transverse mass, defined as $\sqrt{2p_T^\ell p_T^\nu (1 - \cos(\phi^\ell - \phi^\nu))}$ where the measured E_T^{miss} vector provides the neutrino information.

⁴The term ‘fake’ leptons here refers to hadrons mimicking lepton signatures and to leptons arising from heavy hadron decays, whereas ‘real’ leptons come from W and Z decays.

fault muon selection. The fake lepton efficiencies are determined using a low m_T control region $m_T < 20$ GeV with an additional cut $E_T^{\text{miss}} + m_T < 60$ GeV. The efficiencies for signal and fake leptons are parameterised as a function of muon $|\eta|$ and p_T in order to account for the variation of the muon detector acceptance and the profile of hadronic activity in the detector that affects the muon isolation.

For the multijet background estimate in the electron channel, the loose data sample is defined by considering events with electrons passing looser identification criteria. The electron isolation requirement is also modified: the total energy in a cone of $\Delta R = 0.2$ around the electron is required to be smaller than 6 GeV (instead of 3.5 GeV), after correcting for energy deposits from pile-up interactions and for the energy associated with the electron. The fake lepton efficiencies are determined using a low E_T^{miss} control region ($5 \text{ GeV} < E_T^{\text{miss}} < 20 \text{ GeV}$).

In both channels contributions from W +jets and Z +jets backgrounds in the control region, estimated using Monte Carlo simulation, are subtracted.

5.2 W +jets background estimation

At the LHC the rate of W^+ +jets is larger than that of W^- +jets because there are more valence u quarks than d quarks in the proton. Theoretically, the ratio of W^+ +jets and W^- +jets cross sections is predicted much more precisely than the total W +jets cross section [37, 38]. This asymmetry is exploited here to measure the total W +jets background from the data.

Since, to a good approximation, processes other than W +jets give equal numbers of positively and negatively charged leptons, the formula

$$N_{W^+} + N_{W^-} = \left(\frac{r_{MC} + 1}{r_{MC} - 1} \right) (D^+ - D^-), \quad (2)$$

can be used to estimate the total number of W events in the selected sample. Here $D^+(D^-)$ are the total numbers of events in data passing the selection cuts described in Section 4.2 (apart from the b -tagging requirement) with positively (negatively) charged leptons, and $r_{MC} \equiv \frac{N(pp \rightarrow W^+)}{N(pp \rightarrow W^-)}$ is evaluated from Monte Carlo simulation, using the same event selection.

The ratio r_{MC} is found to be 1.56 ± 0.06 in the electron channel and 1.65 ± 0.08 in the muon channel. The dominant uncertainties on r_{MC} originate from those of the parton distribution functions, the jet energy scale, and the heavy flavour fractions in W +jets events (fractions of W +jets events containing $b\bar{b}$ pairs, $c\bar{c}$ pairs and c quarks).

Since the theoretical prediction for heavy flavour fractions in W +jets suffers from large uncertainties, a

data-driven approach was developed to constrain these fractions with some inputs from MC simulation. In this approach samples with a lower jet multiplicity, obtained from the selection described in Section 4.2, but requiring precisely one or two jets instead of four or more jets, are analysed. The numbers $W_{i,\text{pretag}}^{\text{Data}}$, $W_{i,\text{tagged}}^{\text{Data}}$, of $W+i$ jet events in these samples (where $i = 1, 2$), before and after applying the b -tagging requirement, are computed by subtracting the small contributions of other Standard Model processes - electroweak (WW , WZ , ZZ and Z +jets) and top ($t\bar{t}$ and single top) using predictions from the simulation, and by subtracting the multijet background as described in Section 5.1.

A system of two equations, expressing the number of $W+1$ jet events and $W+2$ jets events before and after b -tagging, can be written with six independent flavour fractions as the unknowns, corresponding to fractions of $Wb\bar{b}$ +jets, $Wc\bar{c}$ +jets, and Wc +jets events in the one and two jet bins. The simulation prediction for the ratio of the heavy flavour fractions between the one and two jet bins is used to relate the heavy flavour fractions in the two bins, reducing the number of independent fractions to three. Finally, the ratio of the fractions of $Wc\bar{c}$ +jets and $Wb\bar{b}$ +jets events in the two-jet bin is taken to be fixed to the value obtained from simulated events in order to obtain two equations for two independent fractions. Based on this measurement, the heavy flavour fractions in simulated W +jets events are adjusted by a scale factor 1.63 ± 0.76 for $Wb\bar{b}$ +jets and $Wc\bar{c}$ +jets events and 1.11 ± 0.35 for Wc +jets. When applied to the signal region, an additional 25% uncertainty on these fractions is added, corresponding to the uncertainty of the Monte Carlo prediction for the ratio of heavy flavour fractions in different jet multiplicities. The heavy flavour scale factors are applied to simulated W +jets events throughout this paper, and the effect of their uncertainties on the value of r_{MC} is evaluated.

Using Equation (2), the total number of W +jets events passing the event selection described in Section 4.2 without requiring a b -tagged jet, $W_{\geq 4,\text{pretag}}$, is evaluated to be 5400 ± 800 (stat. + syst.) in the electron channel and 8600 ± 1200 (stat. + syst.) in the muon channel.

The number of W +jets events passing the selection with at least one b -tagged jet is subsequently evaluated as [36]

$$W_{\geq 4,\text{tagged}} = W_{\geq 4,\text{pretag}} \cdot f_{2,\text{tagged}} \cdot k_{2 \rightarrow \geq 4}. \quad (3)$$

Here $f_{2,\text{tagged}} \equiv W_{2,\text{tagged}}^{\text{Data}}/W_{2,\text{pretag}}^{\text{Data}}$ is the fraction of $W+2$ jets events passing the requirement of having at least one b -tagged jet, and $k_{2 \rightarrow \geq 4} \equiv f_{\geq 4,\text{tagged}}^{\text{MC}}/f_{2,\text{tagged}}^{\text{MC}}$ is the ratio of the fractions of simulated W +jets events passing the requirement of at least one b -tagged jet,

for at least four and two jets, respectively. The value of $f_{2,\text{tagged}}$ is found to be 0.065 ± 0.005 in the electron and 0.069 ± 0.005 in the muon channel, where the uncertainties include statistical and systematic contributions. The ratio $k_{2 \rightarrow \geq 4}$ is found to be 2.52 ± 0.36 in the electron channel and 2.35 ± 0.34 in the muon channel. The uncertainties include both systematic contributions and contributions arising from the limited number of simulated events. The total number of W +jets events passing the selection with a b -tagged jet, $W_{\geq 4,\text{tagged}}$, is evaluated to be 880 ± 200 (stat. + syst.) in the electron channel and 1390 ± 310 (stat. + syst.) in the muon channel.

5.3 Other backgrounds

The numbers of background events coming from single top production, Z +jets and diboson events are evaluated using Monte Carlo simulation normalized to the relevant NNLO cross sections for Z +jets events and NLO for diboson events.

5.4 Event yield

The final numbers of expected and observed data events in both channels after the full event selection are listed in Table 1. The number of events in the electron channel is significantly lower than in the muon channel due to the higher lepton p_T requirement and the more stringent missing momentum requirement, which are necessary to reduce the contribution from the multijet background. The overall agreement between expectation and data is good.

6 Reconstruction of the $t\bar{t}$ final state

To measure the charge asymmetry in top pair events, the full $t\bar{t}$ system is reconstructed. For this purpose, a kinematic fit is used that assesses the compatibility of the observed event with the decays of a top-antitop pair based on a likelihood approach.

The likelihood takes as inputs the measured energies, pseudorapidities and azimuthal angles of four jets, the measured energy of the lepton, and the missing transverse momentum. If there are more than four jets in the event satisfying $p_T > 25$ GeV and $|\eta| < 2.5$, all subsets of four jets from the five jets in the event with highest p_T are considered.

The likelihood is computed as

$$L = \mathcal{B}(\tilde{E}_{p,1}, \tilde{E}_{p,2} | m_W, \Gamma_W) \cdot \mathcal{B}(\tilde{E}_{lep}, \tilde{E}_\nu | m_W, \Gamma_W) \cdot$$

Channel	μ + jets pretag	μ + jets tagged	e + jets pretag	e + jets tagged
$t\bar{t}$	7200 \pm 600	6300 \pm 500	4800 \pm 400	4260 \pm 350
W+jets	8600 \pm 1200	1390 \pm 310	5400 \pm 800	880 \pm 200
Single top	460 \pm 40	366 \pm 32	320 \pm 28	256 \pm 22
Z+jets	940 \pm 330	134 \pm 47	760 \pm 270	110 \pm 40
Diboson	134 \pm 7	22 \pm 2	80 \pm 5	13 \pm 1
Multijets	1500 \pm 800	500 \pm 500	900 \pm 500	250 \pm 250
Total background	11700 \pm 1400	2400 \pm 600	7500 \pm 900	1500 \pm 320
Signal + background	18900 \pm 1600	8800 \pm 800	12000 \pm 1000	5800 \pm 500
Observed	19639	9124	12096	5829

Table 1 Numbers of events observed in data and expected from $t\bar{t}$ signal events and various background processes for the pretag and tagged samples defined in Section 4.2. The experimentally determined uncertainties quoted for W +jets and multijet backgrounds include systematic uncertainties on the normalisation. The quoted uncertainties on the other backgrounds are those from theory, taken to be 8% for $t\bar{t}$ and single top, 34% for Z +jets and 5% for diboson backgrounds. The numbers correspond to an integrated luminosity of 1.04 fb^{-1} in both electron and muon channels.

$$\mathcal{B}(\tilde{E}_{p,1}, \tilde{E}_{p,2}, \tilde{E}_{p,3} | m_t, \Gamma_t) \cdot \mathcal{B}(\tilde{E}_{lep}, \tilde{E}_\nu, \tilde{E}_{p,4} | m_t, \Gamma_t) \cdot \mathcal{W}(\hat{E}_x^{\text{miss}} | \tilde{p}_{x,\nu}) \cdot \mathcal{W}(\hat{E}_y^{\text{miss}} | \tilde{p}_{y,\nu}) \cdot \mathcal{W}(\hat{E}_{lep} | \tilde{E}_{lep}) \cdot \prod_{i=1}^4 \mathcal{W}(\hat{E}_{jet,i} | \tilde{E}_{p,i}) \cdot \prod_{i=1}^4 P(\text{tagged} | \text{parton flavour}),$$

where:

- Symbols \mathcal{B} represent Breit-Wigner functions, evaluated using invariant masses of sums of appropriate parton and lepton four-vectors. The pole masses of the W boson and the top quark are fixed to $m_W = 80.4 \text{ GeV}$ and $m_t = 172.5 \text{ GeV}$, respectively. Their widths are taken to be $\Gamma_W = 2.1 \text{ GeV}$ and $\Gamma_t = 1.5 \text{ GeV}$.
- Symbols \mathcal{W} represent the transfer functions associating the reconstructed quantities (\hat{X}) to quarks and leptons produced in the hard scattering (\tilde{X}). $\tilde{E}_{p,i}$ are the energies of partons associated to jets with measured energies $\hat{E}_{jet,i}$. These transfer functions are derived from Monte Carlo simulation.
- $P(\text{tagged} | \text{parton flavour})$ is the b -tagging probability or rejection efficiency, depending on the parton flavour, as obtained from Monte Carlo simulation.

The likelihood is maximised with respect to the energies of the partons, the energy of the charged lepton, and the components of the neutrino three-momentum. The assignment of jets to partons which gives the highest likelihood value is selected. Finally, the sign of the charge of the top quark (or anti-quark) decaying into the lepton is determined from the lepton charge.

The overall efficiency for the reconstruction of the correct event topology is found to be 74% in Monte Carlo simulated $t\bar{t}$ events. Only those events where four jets and a lepton are matched to partonic particles are considered for the efficiency computation.

Distributions of the invariant mass and transverse momentum of the reconstructed top-antitop pair are shown in Fig. 1.

7 Unfolding

The measured distributions of top and anti-top rapidities are distorted by detector effects and an event selection bias. To correct for these distortions the experimental distributions are unfolded to the four-vectors of the top quarks before decay.

The relation between a true distribution T_j (assuming, for simplicity, that there is only one observable of interest) and the reconstructed distribution S_i after detector simulation and event selection can be written:

$$S_i = \sum_j R_{ij} T_j \quad (4)$$

where R_{ij} is the response matrix defined as the probability to observe an event in bin i when it is expected in bin j .

The true distribution T_j can be obtained from the observed distribution S_i by inverting the response matrix. The unfolding problem can similarly be formulated for the case of multiple observables. In this analysis, Bayes' theorem is applied iteratively in order to perform the unfolding [39].

The unfolding is performed using response matrices which account for both detector response and acceptance effects. The response matrices are calculated using Monte Carlo events generated with MC@NLO. The unfolding is done separately, after background subtraction, for the inclusive measured distribution of $\Delta|y|$ (a one-dimensional unfolding problem), and the measured

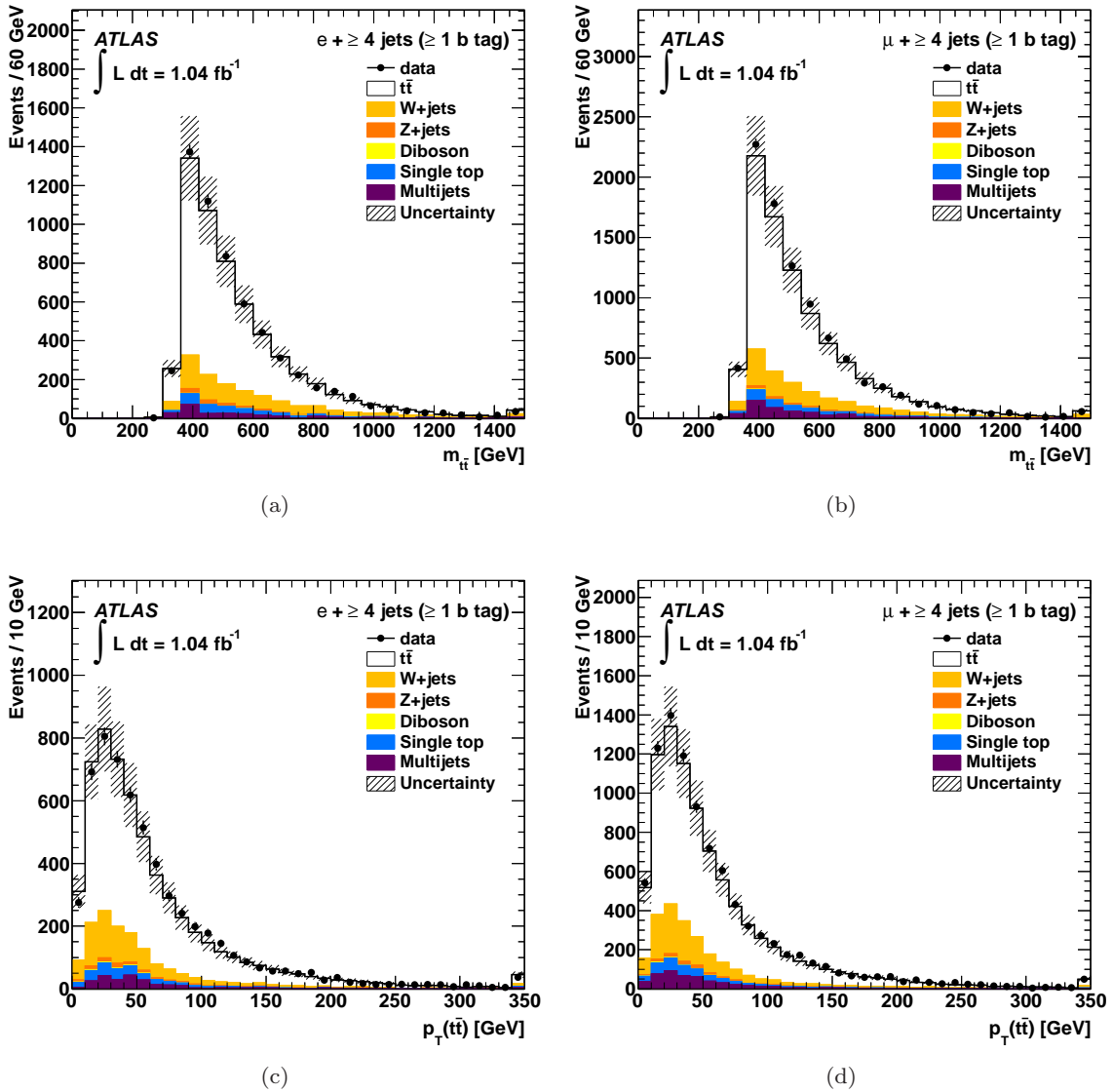


Fig. 1 Expected and observed distributions for the invariant mass (plots (a) and (b)) and transverse momentum (plots (c) and (d)) of the reconstructed $t\bar{t}$ system. The left hand panels show distributions in the electron channel, while the right hand panels show distributions in the muon channel. The data are compared to the sum of the $t\bar{t}$ signal contribution and backgrounds. The background contributions from W +jets and multijet production have been estimated from data, while the other backgrounds are estimated from simulation. The uncertainty on the combined signal and background estimate includes systematic contributions. Overflows are shown in the highest bin of each histogram.

distribution $\Delta|y|$ as a function of the reconstructed top-antitop invariant mass $m_{t\bar{t}}$ (a two-dimensional unfolding problem).

Two bins are used for $m_{t\bar{t}}$ in the two-dimensional unfolding of $\Delta|y|$ versus $m_{t\bar{t}}$, separated at $m_{t\bar{t}} = 450$ GeV. The choice of this $m_{t\bar{t}}$ value is motivated by the observed CDF forward-backward asymmetry [6] and by separating the data sample into two bins with roughly equal number of events.

An additional cut on the value of the likelihood for the $t\bar{t}$ candidate is required in the two-dimensional unfolding, since a large fraction of simulated events with a badly reconstructed $m_{t\bar{t}}$ are found to have a low likelihood value.

The response matrix (including both detector and acceptance effects) for the inclusive A_C measurement is shown in Fig. 2. Six bins in $\Delta|y|$, in the range $-3 < \Delta|y| < 3$, are used in the response matrix, with the outermost bins broader than the inner bins in order

to avoid the occurrence of bins with no entries in the measured distributions. Only a very small fraction of simulated $t\bar{t}$ events are found to have $|\Delta|y| > 3$, and hence such events have a negligible influence on the results.

The unfolding procedure is applied to the observed $\Delta|y|$ distribution in data, after subtracting background contributions. When performing the background subtraction, the shape of the multijet background is obtained by applying the Matrix Method (described in Section 5.1) in bins of $\Delta|y|$. The shape of all remaining backgrounds is taken from Monte Carlo simulation. The value of A_C after unfolding is obtained by counting the numbers of events with $\Delta|y| > 0$ and $\Delta|y| < 0$ in the unfolded $\Delta|y|$ distribution.

8 Systematic uncertainties

Several sources of systematic uncertainties are taken into account in this analysis. These are categorised into the detector modelling, the modelling of signal and background processes and the unfolding method.

8.1 Detector modelling

Small mis-modellings of muon or electron trigger, reconstruction and selection efficiencies in simulation are corrected for by scale factors derived from measurements of the efficiency in data. $Z \rightarrow \mu\mu$ or $Z \rightarrow ee$ and $W \rightarrow e\nu$ decays are used to obtain scale factors as functions of the lepton kinematics. The uncertainties are evaluated by varying the lepton and signal selections and from the uncertainty in the evaluation of the backgrounds. Systematic uncertainties at the level of 1% are found for both cases. The same processes are used to measure the lepton momentum scale and resolution. Scale factors, with uncertainties at the level of (1-1.5)%, are derived to match the simulation to observed distributions. A systematic uncertainty for charge mis-identification of leptons is assigned which is negligible for muons and ranges from 0.2% to 3% for electrons depending on $|\eta|$.

The jet energy scale is derived using information from test-beam data, collision data and simulation. Its uncertainty varies between 2.5% and 8% in the central region, depending on jet p_T and η [33]. This includes uncertainties in the flavour composition of the sample and mis-measurements due to the effect of nearby jets. Pile-up gives additional uncertainties of up to 5% (7%) in the central (forward) region. An extra uncertainty of 0.8% to 2.5%, depending on jet p_T , is assigned to jets arising from the fragmentation of b -quarks, due to differences between light and gluon jets as opposed to jets

containing b -hadrons. The jet energy resolution and reconstruction efficiency are measured in data using techniques described in Refs. [33,40], and their uncertainties are found to be 10% and (1-2)%, respectively.

The b -tagging efficiencies and mis-tag rates are measured in data. Jet p_T dependent scale factors, applied to simulations to match the efficiencies measured in data, have uncertainties which range from 9% to 15% and 11% to 22%, respectively. A systematic uncertainty is assigned for a potential difference of up to 5% between the b -tagging efficiency for b -jets and that of \bar{b} -jets. The uncertainty on the measured luminosity is 3.7% [41,42].

Due to a hardware failure, later repaired, one small region of the liquid argon calorimeter could not be read out in a subset of the data corresponding to 84% of the total integrated luminosity. Data events in which an electron or jet with $p_T > 20$ GeV is close to the affected calorimeter region are rejected for the relevant part of the dataset. Monte Carlo simulated events with electrons or jets of $p_T > 20$ GeV close to the affected region are rejected with a probability equal to the fraction of the integrated luminosity of data for which the calorimeter hardware problem was present. A systematic uncertainty is evaluated by varying the p_T -threshold in data of the electrons and jets near the affected region by ± 4 GeV, corresponding to the uncertainty in the energy lost by objects in the affected region.

8.2 Signal and background modelling

The systematic uncertainty in the modelling of the signal process is assessed by simulations based on different Monte Carlo generators. Sources of systematic uncertainty considered here are the choice of generator and parton shower model, the choice of parton density functions, the assumed top quark mass and the choice of parameters which control the amount of initial and final state radiation. Predictions from the MC@NLO and POWHEG [43,44] generators are compared. The parton showering is tested by comparing two POWHEG samples interfaced to HERWIG and PYTHIA, respectively. The amount of initial and final state radiation is varied by modifying parameters in ACERMC [45] interfaced to PYTHIA according to Ref. [46]. The parameters are varied in a range comparable to those used in the Perugia Soft/Hard tune variations [47]. The impact of the choice of parton density functions is studied using the procedure described in Ref. [48]. MC@NLO samples are generated assuming different top quark masses and their predictions are compared. The observed differences in the results are

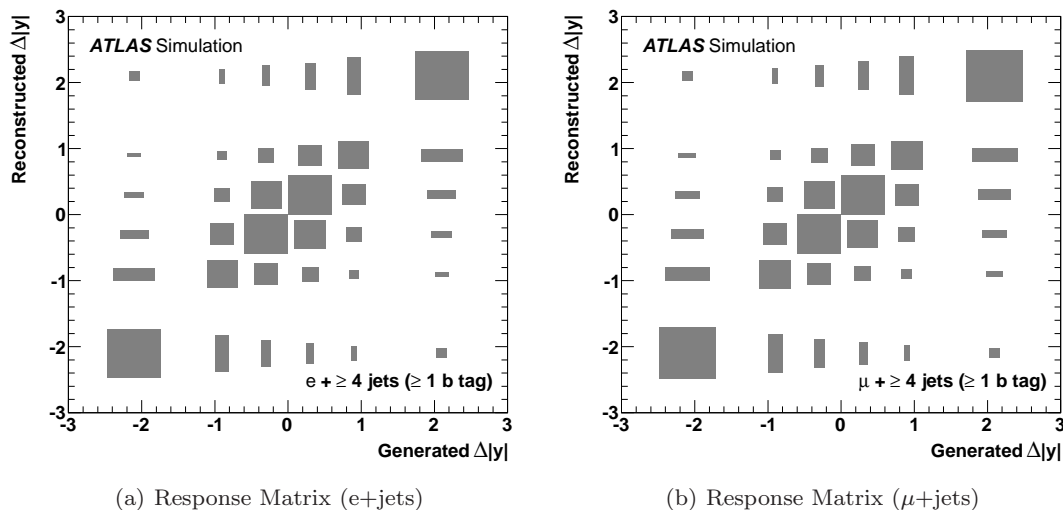


Fig. 2 Correlations between the true and reconstructed values of $\Delta|y|$ encoded in the unfolding response matrix for the electron (left) and muon (right) channels. The value of an entry in the matrix is proportional to the area of the corresponding box.

scaled to variations of ± 0.9 GeV according to the uncertainty on the measured value [49].

As described in Section 5, background processes are either modelled by simulation or estimated in auxiliary measurements. The uncertainty in the estimate of the multijet background is evaluated by considering modified definitions of the loose data sample, taking into account the statistical uncertainty in measurements of $\epsilon_{\text{real}}, \epsilon_{\text{fake}}$ described in Section 5.1 as well as the uncertainties in the normalisations of the W +jets and Z +jets backgrounds which are subtracted in the control region. The total uncertainty is estimated to be 100%. The normalisation of W +jets processes is evaluated from auxiliary measurements using the asymmetric production of positively and negatively charged W bosons in W +jets events. The uncertainty is estimated to be 21% and 23% in the four jet bin, for the electron and muon channels respectively. This uncertainty was estimated by evaluating the effect on both r_{MC} and $k_{2 \rightarrow \geq 4}$ from the JES uncertainty and different PDF and generator choices. Systematic uncertainties on the shape of W +jets distributions are assigned based on differences in simulated events generated with different simulation parameters. Scaling factors correcting the fraction of heavy flavour contributions in simulated W +jets samples are estimated in auxiliary measurements, as described in Section 5.2. The systematic uncertainties are found by changing the normalisations of the non- W processes within their uncertainties when computing $W_{1,\text{pretag}}^{\text{Data}}, W_{1,\text{tagged}}^{\text{Data}}$, as well as taking into account the impact of uncertainties in b -tagging efficien-

cies. The total uncertainties are 47% for $Wb\bar{b}$ +jets and $Wc\bar{c}$ +jets contributions and 32% for Wc +jets contributions. The normalisation of Z +jet events is estimated using Berends-Giele-scaling [50]. The uncertainty in the normalisation is 48% in the four jet bin and increases with the jet multiplicity. A systematic uncertainty in the shape is accounted for by comparing simulated samples generated with ALPGEN and SHERPA [51]. The uncertainty on the normalisation of the small background contributions from single top and diboson production is estimated to be about 10% (depending on the channel) and 5%, respectively.

Limited Monte Carlo sample sizes give rise to a systematic uncertainty in the response matrix. This is accounted for by independently varying the bins of the response matrix according to Poisson distributions.

8.3 Uncertainties from unfolding

Closure tests are performed in order to check the validity of the unfolding procedure. Reweighted $t\bar{t}$ samples with different amounts of asymmetry are considered. Pseudoexperiments are performed, varying the entries in histograms of the reconstructed distribution, to confirm that the response of the unfolding is linear in the true value of A_C and that the true value of A_C is recovered on average. A total of 40 iterations are used in both channels for the inclusive A_C measurement. For the measurement of A_C as a function of $m_{t\bar{t}}$, 80 iterations are used. The number of iterations is chosen by ensuring that the unfolding procedure has converged in

the sense that the absolute change in the unfolded value of A_C after performing an extra iteration is less than 0.001. It is found that the unfolded values of A_C from all pseudoexperiments and the data converge before the chosen numbers of iterations. The potential bias arising from the choice of convergence criterion is taken into account by adding an additional systematic uncertainty corresponding to the change in the unfolded value of A_C obtained by further increasing the number of iterations to very large values (10^5).

Pull distributions are constructed from pseudoexperiments and a relative shift of between 0% and 10% is found in the unfolded value of A_C with respect to the true value. An extra systematic uncertainty is assigned to the unfolded value of A_C obtained from data, corresponding to this shift.

In pseudoexperiments, a small bias is observed in the unfolded distributions corresponding to a relative difference of a few percent between the unfolded result and true value in each bin. An additional relative uncertainty of (2-5)% is applied to all bins of the unfolded distributions, corresponding to the largest relative bin deviation observed in pseudoexperiments.

The statistical uncertainty in the unfolded measurement was computed using pseudoexperiments, propagating the uncertainties from the measured distribution using the statistical correlation matrix.

8.4 Impact of systematic uncertainties

The impact of the systematic uncertainties is evaluated by modifying the subtracted background before unfolding and by modifying the response matrix used for unfolding when relevant. In particular the detector modelling systematic uncertainties are evaluated by shifting the estimated background as well as modifying the response matrix. Signal modelling uncertainties are computed by replacing the response matrix, and background modelling uncertainties by modifying the estimated background.

Table 2 summarises the sources of systematic uncertainties for the inclusive measurement of the charge asymmetry, and their impact on the measured asymmetry, after unfolding. The systematics for the two $m_{t\bar{t}}$ bins are determined in a similar fashion. The evaluation of some systematic uncertainties is limited by the finite size of the Monte Carlo samples. In these cases, the larger of the electron and muon channel uncertainties is used for the uncertainty on the combined result. The resulting combined systematic uncertainties are ± 0.028 in the electron channel and ± 0.024 in the muon channel.

9 Summary of Results

The measured distributions of the top-antitop rapidity difference $\Delta|y| = |y_t| - |y_{\bar{t}}|$ before unfolding are shown in Fig. 3 for the electron and muon channel. Fig. 4 shows the corresponding $\Delta|y|$ distributions after unfolding. After unfolding, the bins of the measured distribution have statistical and systematic correlations. Adjacent bins of the $\Delta|y|$ distributions are found to be statistically anti-correlated with negative correlation coefficients of up to -0.6, whereas other correlations are small.

The measured values of the top charge asymmetry before and after unfolding, defined by Eq. (1) in terms of $\Delta|y|$, are summarised in Table 3. The analytic best linear unbiased estimator (BLUE) method [52, 53] is used to combine the measurement in the electron and muon channels after correction for detector resolution and acceptance.

The measured asymmetries are:

$$A_C = -0.019 \pm 0.028 \text{ (stat.)} \pm 0.024 \text{ (syst.)}$$

for the integrated sample, and

$$A_C = -0.052 \pm 0.070 \text{ (stat.)} \pm 0.054 \text{ (syst.)}$$

for $m_{t\bar{t}} < 450$ GeV,

$$A_C = -0.008 \pm 0.035 \text{ (stat.)} \pm 0.032 \text{ (syst.)}$$

for $m_{t\bar{t}} > 450$ GeV.

The measurement for the integrated sample can be compared with the result of the CMS Collaboration, $A_C = -0.013 \pm 0.028 \text{ (stat.)}^{+0.029}_{-0.031} \text{ (syst)}$ [19]. Fig. 5 summarizes the measurements for the two $m_{t\bar{t}}$ regions. These results are compatible with the prediction from the MC@NLO Monte Carlo generator of $A_C = 0.006 \pm 0.002^5$, showing no evidence for an enhancement from physics beyond the Standard Model.

10 Comparison of LHC and Tevatron Results

The measurement of the charge asymmetry at the LHC is a test of the unexpectedly large forward-backward asymmetry observed at the Tevatron. However, because the LHC is a pp collider and the centre of mass energy is around three times larger, any relation between the two asymmetries is model-dependent. Here a comparison is made between the predicted values of the Tevatron and

⁵The prediction of 0.0115 ± 0.0006 for the charge asymmetry found in Ref. [54] differs from the MC@NLO prediction of 0.006 ± 0.002 , due to the former taking the LO prediction for the denominator in the definition (1) of A_C , and taking into account QED effects. The uncertainty on the MC@NLO prediction is obtained by considering variations in the renormalisation and factorisation scales and different sets of PDFs.

Source of systematic uncertainty on A_C	Electron channel	Muon channel
<i>Detector modelling</i>		
Jet energy scale	0.012	0.006
Jet efficiency and resolution	0.001	0.007
Muon efficiency and resolution	<0.001	0.001
Electron efficiency and resolution	0.003	0.001
b-tag scale factors	0.004	0.002
Calorimeter readout	0.001	0.004
Charge mis-ID	<0.001	<0.001
b-tag charge	0.001	0.001
<i>Signal and background modelling</i>		
Parton shower/fragmentation	0.010	0.010
Top mass	0.007	0.007
$t\bar{t}$ modelling	0.011	0.011
ISR and FSR	0.010	0.010
PDF	<0.001	<0.001
W+jets normalization and shape	0.008	0.005
Z+jets normalization and shape	0.005	0.001
Multijet background	0.011	0.001
Single top	<0.001	<0.001
Diboson	<0.001	<0.001
MC Statistics	0.006	0.005
Unfolding convergence	0.005	0.007
Unfolding bias	0.004	<0.001
Luminosity	0.001	0.001
Total systematic uncertainty	0.028	0.024

Table 2 List of sources of systematic uncertainties and their impact on the measured asymmetry in the electron and muon channel. In cases where asymmetric uncertainties were obtained, a symmetrisation of the uncertainties was performed by taking the average of the absolute deviations under systematic shifts from the nominal value.

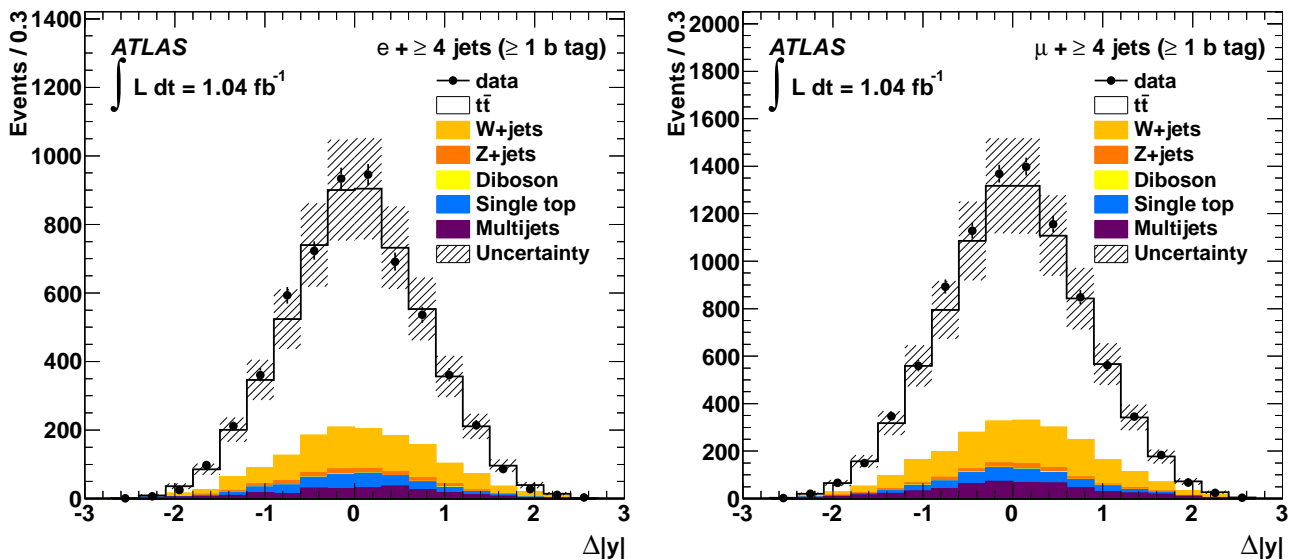


Fig. 3 The measured $\Delta|y|$ distribution before unfolding for the electron channel (left) and for the muon channel (right) after b -tagging is applied. Data (points) and Monte Carlo estimates (solid lines) are represented. The multijet background and the normalisation of the W +jets background are obtained as explained in Section 5. The uncertainty on the combined signal and background estimate includes both statistical and systematic contributions.

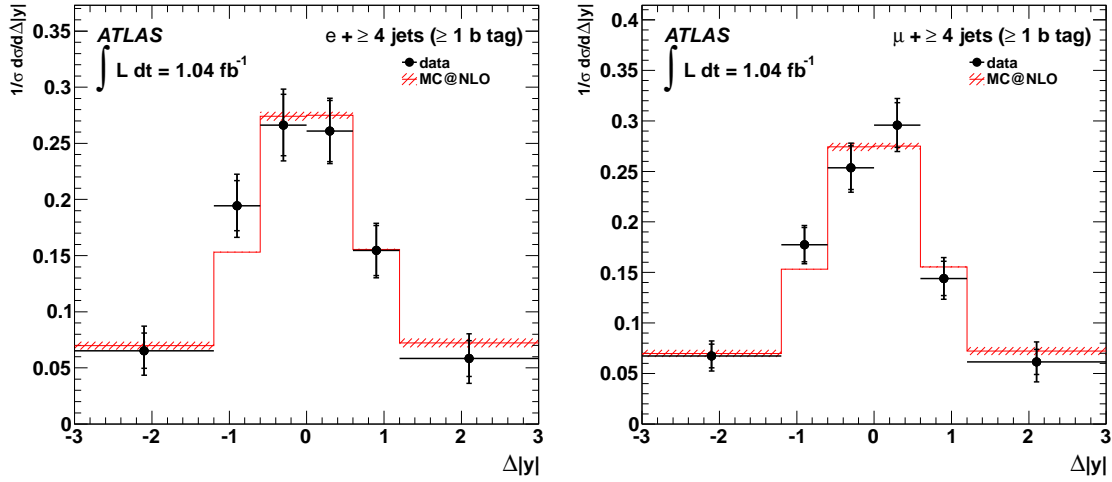


Fig. 4 The unfolded $\Delta|y|$ distribution for the electron channel (left) and the muon channel (right) after b -tagging, compared to the prediction from MC@NLO. The uncertainties on the measurement include both statistical and systematic contributions, which are shown separately. The inner part of the error bars corresponds to the statistical component of the uncertainty, while the outer part corresponds to the systematic component. The error bands on the MC@NLO prediction include uncertainties from parton distribution functions and renormalisation and factorisation scales.

Asymmetry	reconstructed	detector and acceptance unfolded
A_C (electron)	-0.034 ± 0.019 (stat.) ± 0.010 (syst.)	-0.047 ± 0.045 (stat.) ± 0.028 (syst.)
A_C (muon)	-0.010 ± 0.015 (stat.) ± 0.008 (syst.)	-0.002 ± 0.036 (stat.) ± 0.024 (syst.)
Combined		-0.019 ± 0.028 (stat.) ± 0.024 (syst.)

Table 3 The measured inclusive charge asymmetry values for the electron and muon channels after background subtraction, before and after unfolding.

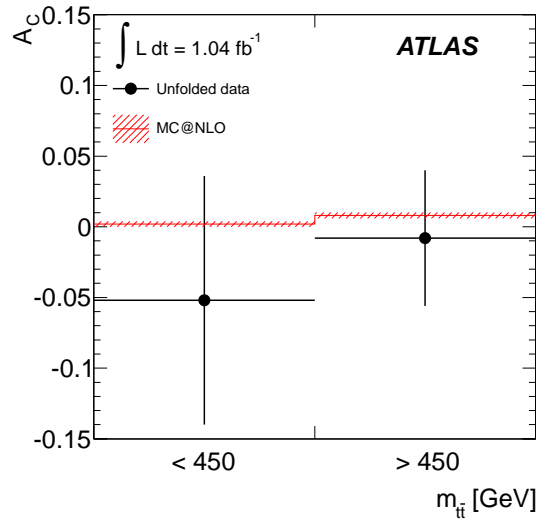


Fig. 5 Unfolded asymmetries in two regions of $m_{t\bar{t}}$ compared to the prediction from MC@NLO. The error bands on the MC@NLO prediction include uncertainties from parton distribution functions and renormalisation and factorisation scales.

LHC asymmetries for a few simple models beyond the SM. These are: (i) a flavour-changing Z' boson with right-handed couplings, exchanged in the t channel in $u\bar{u} \rightarrow t\bar{t}$ [10]; (ii) a W' boson, also with right-handed

couplings, contributing in $d\bar{d} \rightarrow t\bar{t}$ [11]; a heavy axigluon \mathcal{G}_μ exchanged in the s channel [8,9]; (iv) a scalar doublet ϕ , with the same quantum numbers as the SM Higgs [55]; (v) a charge 4/3 scalar, colour-sextet (Ω^4)

or colour-triplet (ω^4), contributing in the u channel to $u\bar{u} \rightarrow t\bar{t}$ [12, 13]. In all these models, the parameter space is described by the mass M of the new particle (except for the axigluon which is assumed heavy, with $M \gg 7$ TeV) and a single coupling g .

In order to find the correlated predictions for the forward-backward and charge asymmetries in each model, a comprehensive scan over the mass M and the coupling g is performed using the PROTOS generator [56], considering masses between 100 GeV and 10 TeV and the range of couplings for which the new physics contribution to the $t\bar{t}$ cross section at the Tevatron lies in the interval $[-0.8, 1.7]$ pb. This is a conservative requirement which takes into account the different predictions for the SM cross section as well as the experimental measurement (see Ref. [17] for details).

In addition, a conservative upper limit on new physics contributions to $\sigma_{t\bar{t}}$ for $m_{t\bar{t}} > 1$ TeV is imposed. Further details can be found in Refs [17, 55]. The coloured areas in Fig. 6 (a) represent the ranges of predicted values for the inclusive Tevatron forward-backward asymmetry, A_{FB} , and the inclusive LHC charge asymmetry, A_C , for the new physics models. The new physics contributions are computed using the tree-level SM amplitude plus the one(s) from the new particle(s). To a good approximation, the total asymmetries A_{FB} , A_C are obtained from the former by summing the SM contribution (at NLO in the lowest order). The horizontal lines correspond to the present ATLAS measurement and the measurement reported by the CMS Collaboration [19]. The vertical lines correspond to the asymmetry measurements at the Tevatron, $A_{FB} = 0.158 \pm 0.075$ [6] and $A_{FB} = 0.196 \pm 0.065$ [7].

The ATLAS charge asymmetry measurement disfavours models with a new flavour-changing Z' or W' vector boson proposed to explain the measured Tevatron asymmetry. Minimal Z' models are also excluded by the non-observation of same-sign top quark production [57]. For the other new physics models the asymmetries measured at the Tevatron are consistent with this measurement, within the experimental uncertainties.

Fig. 6 (b) shows the allowed regions for the high-mass asymmetries ($m_{t\bar{t}} > 450$ GeV) at the Tevatron and the LHC for the six new physics models. The vertical lines represent the CDF measurement $A_{FB} = 0.475 \pm 0.114$ [6], while the horizontal lines correspond to the present ATLAS measurement. In both panels of Fig. 6, the range of variation of SM predictions found in Refs [54, 58, 59] is indicated by a box. The predictions of the six new physics models are in tension with the CDF and ATLAS high-mass measurements considered together.

11 Conclusion

To summarize, the top quark charge asymmetry was measured in $t\bar{t}$ events with a single lepton (electron or muon), at least four jets and large missing transverse momentum using an integrated luminosity of 1.04 fb^{-1} recorded by the ATLAS experiment at a centre of mass energy of $\sqrt{s} = 7$ TeV. The reconstruction of $t\bar{t}$ events was performed using a kinematic fit. The reconstructed inclusive distribution of $\Delta|y|$ and the distribution as a function of $m_{t\bar{t}}$ were unfolded after background subtraction to obtain results that can be directly compared with theoretical computations. The results are compatible with the prediction from the MC@NLO Monte Carlo generator. These measurements disfavour models with a new flavour-changing Z' or W' vector boson that have been suggested to explain the measured Tevatron asymmetry.

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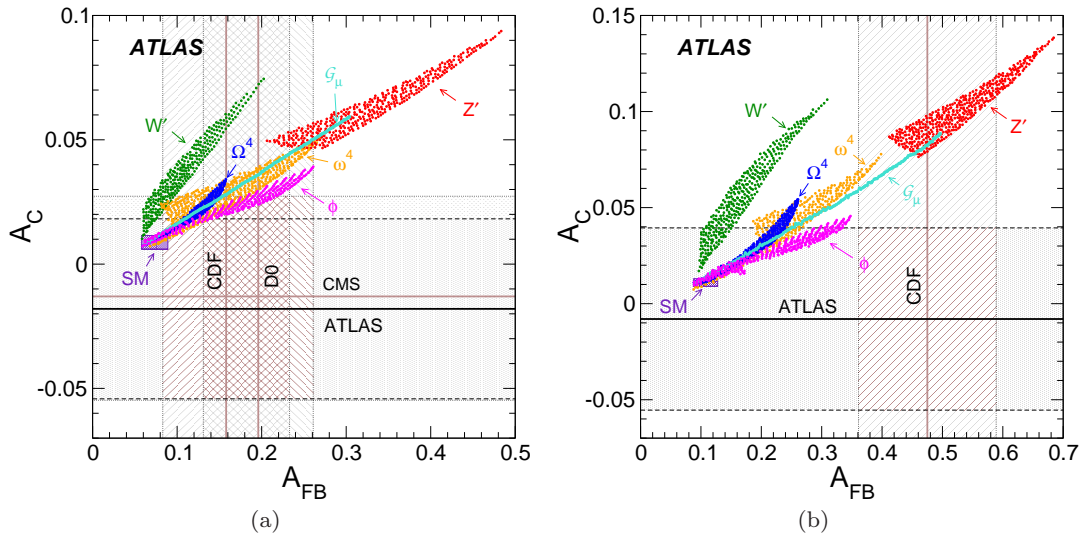


Fig. 6 Measured FB asymmetries from the Tevatron and charge asymmetries from the LHC, compared to predictions from the SM as well as predictions incorporating various potential new physics contributions. The horizontal (vertical) bands and lines correspond to the ATLAS and CMS (CDF and D0) measurements. In (a) the inclusive values are presented and in (b) the ATLAS measurement for $m_{t\bar{t}} > 450$ GeV is compared to the CDF measurement. The MC predictions for the new physics models are from Refs. [17, 55].

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